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Ligand Hyperfine Interactions in Oxomolybdenum(V) and Oxochromium(V) Halide Complexes

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Ligand superhyperfine interactions have been observed in esr spectra of solutions, frozen glasses, and diamagnetically diluted single crystals of $MOBr_3^{2-}$, $MOBr_3^{2-}$, and $CroCl_3^{2-}$. The results indicate considerable covalency in metal-ligand bonding. In the molybdenum and chromium complexes the superhyperfine tensors indicate π bonding involving the metal d_{xy} orbital and p orbitals of the four equatorial ligands. In these d' complexes the ligand orbital involved in bonding possesses some s character.

Ligand superhyperfine interaction has been observed for fluoro-,¹⁻⁴ chloro-,^{1,5} and bromooxomolybdate(V)^{6,7} complexes. The presence of isotropic fluorine hyperfine splittings in a solution of $CrO₃$ dissolved in HF has been reported by Garifyanov⁸ but the spectra were not analyzed in detail. Verbeek and Cornaz⁹ have analyzed the spectrum of $CrOF_b^{2-}$ in 38% HF in the glassy state, by computer simulation, but values for ligand hyperfine splitting tensors were given by Manoharan and Rogers,¹ who also made single-crystal esr measurements.1 **A** detailed single-crystal esr investigation of $MoOCl₅²⁻ has been made and the theory was given.¹⁰$ Similarly a careful esr study of $CrOCl₅²⁻$ has been carried out by Kon and Sharpless.¹¹

We have attempted to perform a complete analysis of ligand superhyperfine interactions in oxomolybdate(V) and $\operatorname{oxochromatic}(V)$ halide complexes by studying the esr spectra under a wide range of conditions. An examination of single-crystal spectra is the most accurate method of obtaining all of the elements of the ligand superhyperfine tensors, and particular attention was paid to such studies.

Experimental Section

Preparation of Complexes.--Complexes of $Mo(V)$ of the type M_2MoOX_{δ} , where $M^+ = K^+$, NH_4^+ and $X^- = Br^-$, I⁻, were prepared by the method of Allen, et $al.^{12}$ K₂CrOCl₅ was also made according to Weinland and Friederer.I3

Growth of Single Crystals.--($NH_4)_2MoORr_5$ in $(NH_4)_2InBr_5$. HzO was prepared by the method of Wentworth and Piper.I4 Single crystals of $(NH_4)_2$ CrOCl₅ in $(NH_4)_2$ InCl₆·H₂O were pre-

- **(3)** J. **T. C. Van Kemenade,** J. **L. Verbeek, and P.** F. **Cornaz,** *Rec. Trav.* **A.** N. **Ermakov,** *Dokl. Akad. Nauk SSSR,* **169, 1107 (1966).** *Chim.,* **86, 629 (1966).**
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- **(11) H. Kon and** N. **E. Sharpless,** *ibid.,* **49, 006 (1965).**
- **(12) E. A. Allen,** B. J. **Brisdon,** D. **A, Edwards,** *G.* **W. A. Fowles, and R.** *G.* **Williams,** *J. Chem. Soc.,* **4649 (1963).**
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pared in the manner used for the oxomolybdate(V)-chloro complex.

Instruments.-Spectra were taken at 9.2 , 9.5 , 35.0 , and 35.5 GHz by using Varian X-band and Q-band esr spectrometers. Spectra were recorded at temperatures of **4.2,** *77,* and 300'K. Temperatures intermediate between *77* and 300'K were investigated by employing a Varian variable-temperature control apparatus with liquid nitrogen as coolant.

Particular care was taken to avoid saturation of the hyperfine and superhyperfine resonances at low temperature. Modulational amplitudes were kept to a minimum.

Diphenylpicrylhydrazyl (DPPH), pitch in potassium chloride, and peroxylaminedisulfonate were used as reference standards to facilitate the calculation of **g** values.15

Analysis of the Spectra.--A computer program was written to synthesize the spectra of crystals of rhombic symmetry. This program is similar to the spectrum generation program developed by Venable.¹⁶

The esr spectra of magnetically dilute powders, prepared from powdering single crystals, were simulated by summing spectra corresponding to particular orientations over all angles. Angular increments of approximately 1° were taken. We were not successful in simulating ligand hyperfine splittings, but since spin densities, calculated from metal and ligand hyperfine spectra agree well, we have confidence in the ligand hyperfine constants.

Results

 $MoOBr₅²$. We were unable to observe bromine superhyperfine interaction in solutions of the $MoOBr₅²$ complex; however, we were able to observe bromine superhyperfine structure for $MoOBr₅²⁻$ in frozen acid glasses. The spectra did not appear to be particularly sensitive to the concentration of $MoOBr₅²⁻$ or the acidity of the solution used to prepare the glass. **A** typical spectrum showed no detectable superhyperfine interaction of the molybdenum parallel hyperfine components on the g_{\parallel} absorption; however, the g_{\perp} absorption displayed a bromine superhyperfine structure consisting of 13 nearly equally spaced lines (see Figure 1). Dowsing and Gibson' suggested that the paramagnetic electron interacts with four equivalent bromines. However, in agreement with the interpretation of Manoharan and Rogers¹ for $MoOCl₅²⁻$ and $MoOF₆²⁻, we believe the hyperfine interaction arises$ from two sets of two equivalent bromines. The

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Figure 1.-Spectrum of $MoOBr₆²⁻$ in frozen acid glass.

13 lines arise from the near degeneracy of the molybdenum perpendicular hyperfine components and the bromine superhyperfine components. The values of the magnetic tensor elements measured from solution and frozen glass spectra are given in Tables I and 11.

TABLE I STUDIED (METAL NUCLEAR HYPERFINE CONSTANTS) ESR SPECTRA AND DERIVED QUANTITIES FOR THE COMPLEXES[®] Complex $\langle g \rangle$ $g \parallel$ $g \perp$ $\langle a \rangle$ $A \parallel$ $A \perp$
(NH₄)₂[M₀OB_{r₃] +1.993 +2.090 +1.945 +41.7 +66.0 +30.0} in $(NH_4)_2$ [In- $Br_1 \cdot H_2 O$ (single crystal) in $(NH_4)_2$ [In- $1 \text{ r} \cdot \text{H}_2$ O] (powder)
 $MoOHr₆² in$ **solution** and frozen glasses solution $(N_4)_2[M_0Br_5]$ $+2.090 +1.945$ $+1.994 + 2.09 + 1.944 + 41.5 + 66.0 + 30.0$ $MoOI₅²–in +2.508$ $CrOCl₃² + 1.9877 +18.2$ so ution $CrOCl₃² - (single +2.008 +1.977 +36.1 +9.7)$ crystal)
CrOCl₅² – (frozen

 $+2.008$ $+1.977$ $+36.1$ $+9.7$

 $MoOI₅²$ -.-Although ligand superhyperfine interaction was observed for the $MoOI₅²⁻$ complex in acid solution, the separation of the molybdenum hyperfine components was smaller than for other oxomolybdate-

^{*a} A* values are given in cm⁻¹ \times 10⁻⁴.</sup>

glass)

(V) halide complexes and resulted in considerable overlap of the superhyperfine components, thus making quantitative measurements of either the molybdenum hyperfine interaction or iodine superhyperfine interaction impossible at present.

Crystal and Polycrystalline Spectra.-The Zeeman and metal hyperfine interaction tensors can be thought of as having their origin at the molybdenum nucleus. To calculate the theoretical expressions for the elements of the Zeeman and hyperfine interaction tensors from molecular orbital theory it is convenient to take the \bar{z} axis along the M= \bar{O} bond. The x and y axes are located in the plane formed by the four equatorial halide ligands. This coordinate system permits unambiguous specification of Zeeman and metal hyperfine tensor elements but not the ligand superhyperfine tensor elements.

To consider ligand superhyperfine interaction a eoordinate system which has its origin at the ligand

nucleus rather than the metal nucleus is considered. **A** local coordinate system is chosen on an equatorial halide¹⁷ with the x axis along the Mo-X bond, z parallel to the symmetry axis of the complex (the $M=O$ bond), and *y* chosen to form a right-handed coordinate system.

The electronic Zeeman and metal hyperfine tensors obey the relationships

$$
g = (g_z^2 \cos^2 \theta + g_x^2 \sin^2 \theta \cos^2 \phi + g_y^2 \sin^2 \phi)^{1/z} \quad (1)
$$

 gA (metal) = $(A_z^2g_z^2 \cos^2 \theta + A_x^2g_x^2 \sin^2 \theta \cos^2 \phi +$

$$
A_y^2 g_y^2 \sin^2 \theta \sin^2 \phi
$$
^{1/z} (2)

The spherical polar angles θ and ϕ relate the external magnetic field vector \vec{H} to the *z* and *y* axes, respectively. The orientation for which $\theta = 0$ gives rise to the spectrum corresponding to g_z and A_z while those for which θ $= \pi/2$ and $\phi = 0$ give rise to the spectra corresponding to (g_x, A_x) and (g_y, A_y) , Table I.

 MoORr_5^- .---For $(\text{NH}_4)_2\text{MoORr}_5$ in $(\text{NH}_4)_2\text{InBr}_5\text{H}_2\text{O}$, as for the chloride complex,^{1} ligand superhyperfine structure was not observed at room temperature and the best resolution of the bromine superhyperfine structure was obtained near 77°K. The spectra observed for single crystals of the bromide complex in $(NH_4)_2InBr_5$. H_2O (Figure 2) were very similar to those observed¹ for the chloride complex.

At $\theta = 0$ no bromine superhyperfine structure was observed which indicates $A_x(\text{Br}) \cong 0 \text{ G}$. At $\theta = \pi/2$ and $\phi = 0$ the spectrum consisted of 7 equally spaced lines while at $\theta = \pi/2$ and $\phi = \pi/4$ the spectrum consisted of 13 equally spaced lines and, together with the intensity ratios and the observed splittings, indicates that $A_z(\text{Br}) \cong 0 \text{ G}$, $A_y(\text{Br}) = 40 \pm 1 \text{ G}$, and $A_r(\text{Br}) = 0$ G.

Pentachloro(oxo)chromate(V).—The esr spectrum of K_2 ⁵³CrOCl₅ in acetic acid has been studied by Kon and Sharpless.¹¹ In our experiments $CroCl₅²⁻$ was doped in single crystals of $(NH_4)_2InCl_5 \tcdot H_2O$. The crystals were grown in aqueous HC1. The esr parameters obtained from measurements at low temperatures on these single crystals and powders, obtained by crushing the crystals, are shown in Tables I and 11. Superhyperfine interactions due to 35Cl and 37Cl were observed.

Discussion

Ligand Superhyperfine Interactions.—If we assume that electron-nuclear dipolar interaction accounts for the anisotropic parts of the ligand superhyperfine interaction tensors, the anisotropic part of the superhyperfine interaction tensor may be expressed as

$$
A_{\rm p} = \sqrt[2]{8g_{\perp}\beta_{\rm e}}g_{\rm n}\beta_{\rm n}\langle r^{-3}\rangle^{'}\rho
$$

and the spin density in the halide orbitals may be calculated (Table 111), where they are compared with the results of Manoharan and Rogers.¹

⁽¹⁷⁾ We restrict ourselves **at** this time to considering only superhyperfiue interaction involving the four equatorial **ligands.**

TABLE I1

a From ref 1. A_x , A_y , A_z , and A_z are ligand hyperfine isotropic and isotropic splittings. A_{xz} , A_{yy} , and A_{zz} are calculated by subtracting the isotropic part from the ligand hyperfine constants A_x , A_y , and A_x . A_{pq} and $A_{p\pi}$ have been calculated from corrected dipolar interactions of the unpaired electrons in the in-plane p_x and p_y orbitals of the halide ligands.¹

Spin density in the ligand p_{ν} orbital can most logically be thought of as arising from in-plane π bonding involving the molybdenum b_2 (primarily d_{xy}) and ligand b_2 (primarily p_v) orbitals. Since the paramagnetic electron would be thought to be localized primarily in the ground-state $b_2(d_{xy})$ orbital, it is reasonable that a fairly large paramagnetic density would be expected to reach the ligand $b_2(p_y)$ orbital.

The explanation of spin density in the ligand p_x orbital is not so straightforward. The ground-state b_2 orbital is of the wrong symmetry to overlap appreciably with the ligand p_x orbitals and configurational interaction involving an excited molybdenum $b_1(d_{z^2-y^2})$ orbital must be considered. **u** bonding would be considerably less effective than in-plane *n* bonding involving the metal and ligand b_2 orbitals and it would be less effective for the chloride and bromide ligands than for the fluoride' ligands. There is not very much evidence for any significant amount of out-of-plane π bonding.

Figure 2.—Polycrystalline spectra of $(NH_4)_2MOOR_5$ in $(NH_4)_2$ -
InBr₆·H₂O.

^{*a*} The values of f_{π} and f_{σ} were calculated by using the measured values of A_p and the values of $\langle r_{np}^{-3} \rangle$ tabulated by R. G. Barnes and W. V. Smith, *Phys. Rev.,* **93,** 95 (1954); A. J. Freeman and R. E. Watson, *ibid.,* **123,** 521 (1961); and R. Bersohn and R. G. Shulman, *J. Chem. Phys.*, 45, 2298 (1966). The values of f_s were calculated by using the observed values of *A,* and the atomic or ionic values of $\psi_{ns}(0)$ tabulated by G. W. Chantry, A. Horsfield, J. R. Morton, J. R. Rowlands, and D. H. Whiffen, *Mol. Phys., 5,* 233 (1962); E. Clementi, C. C. J. Roothaan, and M. Yoshimine, *Phys. Rev.,* **127,** 1618 (1962); R. J. Cook, J. R. Rowlands, and D. H. Whiffen, *Proc. Chew. Soc.,* 252, (1962); R. E. Watson and A. J. Freeman, *Phys. Rev.,* **123,** 521 (1961); R. Bersolm and R. G. Shulman, *J. Chem. Phys.,* **45,** 2298 (1966); G. L. Gardner, *ibid.*, 46, 2991 (1967). *b* From ref 1.

The results appear to be consistent with this analysis. Metal-ligand overlap integrals of 0.01-0.20 could easily account for the observed distribution of electron density.

Measurement of the spectra of single crystals and of polycrystalline powders has greatly facilitated the interpretation of the frozen glass spectra of $MoOBr₅²⁻$. Results are listed in Tables 1-111, and the following quantities were calculated as in our previous work:15 $-K = 41.0 \times 10^{-4}$ cm⁻¹, $-\chi = 5.24$ au, $\rho = 48.9 \times$ 10^{-4} cm⁻¹, $\beta^2 = 0.890$

As can be seen in the present results and those of Manoharan and Rogers, $¹$ the large changes in the values</sup> of the Zeeman and metal hyperfine interaction tensors with changes in halide ligand do not apparently arise from changes in metal-ligand bonding and probably arise from changes in the exchange polarization of the ground-state d_{xy} and inner s orbitals of molybdenum. The theory of exchange polarization applicable to m olybdenum $d¹$ complexes has recently been extensively reviewed.¹⁸

We have also employed the molecular orbital treatment of Manoharan and Rogers¹ and calculated the bonding coefficients for $MoOBr₆²⁻$ and $CrOCl₆²⁻$ by use of the quantities given in Table IV and the coefficients

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⁴ Reference 2. ^b Reference 1. ^o T. M. Dunn, *Trans. Faraday* Soc., 57, 1441 (1961). d Reference 6. e Reference 11. f Reference 18.

are $\beta_2 = 0.932$, $\epsilon = 0.844$, and $\beta_1 = 0.596$ for MoOBr₅²⁻¹ and 0.886, 0.958, and 0.952 for $CrOCl₅²$. Although no trend is noted in the coefficients for $\beta_2(d_{xy})$, regular decreases in ϵ and β_1 are recorded in passing from fluoride^t to chloride¹ to bromide.

By use of Mulliken population analysis one can then calculate spin densities and we find 3.3% for equatorial halide in MoOCl₅²⁻ and 5.3% for CrOCl₅²⁻.

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The Preparation and Optical and Electron Spin Resonance Spectra of Some Hexachloro- and Pentachloroalkoxovanadates(IV)

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Several compounds of the type $[R_4N]_2V(OR')Cl_5$ have been prepared from $[(CH_3)_4N]_2VCl_6$ and $[(CH_3CH_2)_4N]_2VCl_6$. All compounds show normal, spin-only paramagnetism for a $d¹$ ion at room temperature. The magnetic moments for the alkoxo complexes vary little with temperature. Electronic, infrared, and electron spin resonance spectra have been measured. Bonding coefficients for the pertinent molecular orbitals have been calculated. $N_{\pi_2}^2$ and $N_{\sigma_2}^2$ decrease with increasing length of alkoxo chain while $N_{\pi_1}^2$ increases. Spin densities have been calculated for the equatorial Cl atoms from the bonding parameters.

Much recent interest has centered on the investigation of $d¹$ transition elements involving oxy ions of the type MO^{n+} where $M^{(n+2)+}$ is Cr^{5+} , Mo^{5+} , or V^{4+} . The halide complexes of these ions have been studied carefully and it is well established that there is substantial π bonding between the metal and oxygen atom.¹⁻⁵ Recent work in this laboratory has concerned the preparation and characterization of d¹ transition element alkoxides and complexes containing alkoxides as ligands.⁶⁻¹² It was decided that vanadium(IV) chloride

(6) R. A. D. Wentworth and C. H. Brubaker, Jr., Inorg. Chem., 2, 551 $(1963).$

alkoxides would be of interest because of the large amount of work which has been done with the vanadyl systems. The monoalkoxide species ought to be sixcoordinate anions where the metal to oxygen bond is formally a single bond. A study of these compounds should show the effect on the bonding parameters of changing the donating strength of the oxygen by varying the alkoxide group. The effect of changing the halogen from F^- to Cl^- to Br^- has been shown in the case of $MoOX₅²–.1.2$

Bradley and Mehta¹³ and Thomas¹⁴ have prepared a number of vanadium alkoxides of the type $V(OR)_4$. Various investigators have isolated substances of the general form $V(OR)_2Cl_2 \cdot ROH$ from alcoholic solutions.¹⁵⁻¹⁷ Several salts of the hexachlorovanadate ion

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